

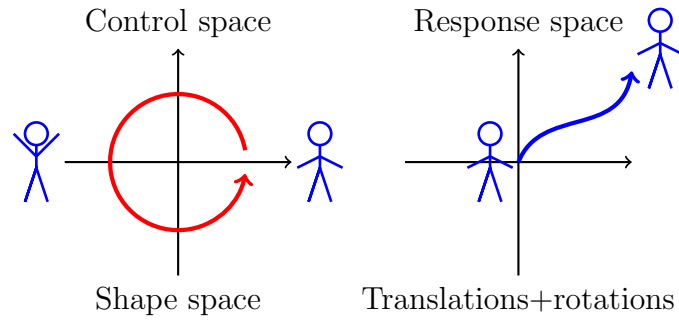
# Adiabatic Quantum Transport II

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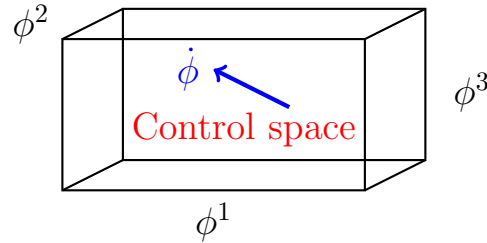
July 12, 2016

# 1 Controls and Response



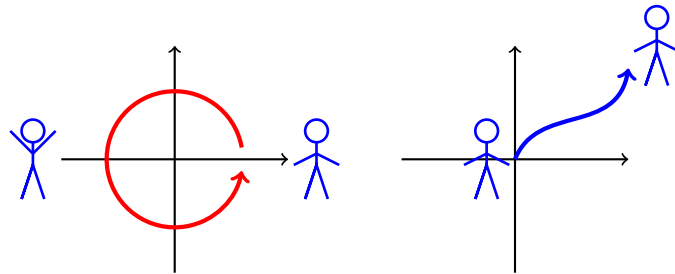
## 1.1 Driving and response

- Controls:  $\phi$
- Hamiltonian  $H(\phi)$
- Driving:  $\underbrace{\dot{\phi} = \frac{d\phi}{dt}}_{\text{vector}}$
- Response:  $\underbrace{d(\text{virtual work})}_{\text{1-form}} = -dH = -\frac{\partial H}{\partial \phi} d\phi$
- Control spaces with interesting topology
- $H$  is N-particle Hamiltonian.



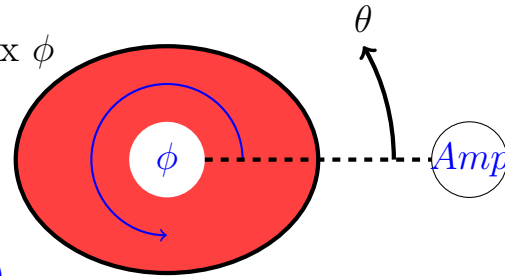
## 1.2 Linear response

- Transport coefficients  $(\partial_j H) = R_{jk} \dot{\phi}^k$
- Dissipation (power):  $(\partial_j H) \dot{\phi}^j$
- Non-dissipative transport:  $(\partial_j H) \dot{\phi}^j = 0 \iff R_{jk} = -R_{kj}$
- Total response:  $\int (\partial_j H) dt = \underbrace{\mathbb{E}(R_{jk})}_{\text{motion per stroke}} \delta\phi^k$
- Compute  $R$  from QM.



### 1.3 Driving loop current

- Control: Aharonov-Bohm flux  $\phi$
- $\mathbf{A} = \frac{\phi}{2\pi} d\theta = \phi \delta(\theta) d\theta$
- Drive: emf  $\dot{\phi} = - \oint E_j dx^j$
- Minimal coupling:  $H(\mathbf{p} - \mathbf{A})$
- Heisenberg equation:  $\mathbf{v} = i[H, \mathbf{x}] = \partial_p H$
- Response : Loop current



$$\begin{aligned}
 -dH &= -(\partial_\phi H) d\phi \\
 &= \frac{1}{2}(\mathbf{v} \cdot \partial_\phi \mathbf{A} + \partial_\phi \mathbf{A} \cdot \mathbf{v}) d\phi \\
 &= \frac{1}{2}(v_\theta \delta(\theta) + \delta(\theta) v_\theta) d\phi
 \end{aligned}$$

- Choosing  $A$  is choosing a cross section.

## 1.4 Electric field driving current

- Bloch Hamiltonians: Hofstadter model

- $H(\mathbf{k}) = e^{ik_1}\mathbf{T}_1 + e^{ik_2}\mathbf{T}_2 + h.c.$

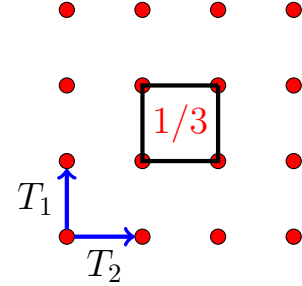
- $\mathbf{T}_1 = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}, \quad \mathbf{T}_2 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^2 \end{pmatrix}$

- $B = 1/3 \Leftrightarrow \omega = e^{2\pi i/3}$

- Control: Brillouin zone  $\mathbf{k} \in T^2$

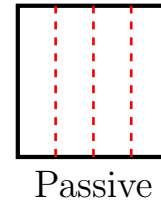
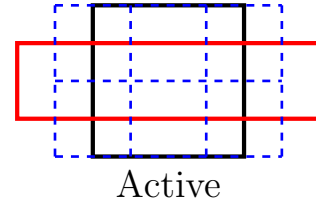
- Driving: External electric field:  $\dot{\mathbf{k}} = \mathbf{E}$

- Response:  $-dH = -\underbrace{(\partial_k H)}_{\text{velocity}} dk = -i(T_j - T_j^*) dk^j$



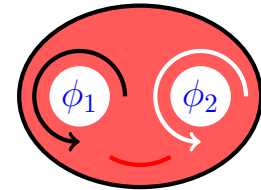
## 1.5 Controlling the metric

- Control: Metric tensor  $g$
- Linear deformations  $L$
- Active :  $\mathbf{x}' = L^{-1}\mathbf{x}$ ; Passive  $g' = L^t g L$
- $(d\ell)^2 = (d\mathbf{x}'^t)g'(d\mathbf{x}') = (d\mathbf{x}^t)g(d\mathbf{x})$
- Driving: strain rate =  $\dot{g}$
- Response: Stress  $dH(g) = -(\partial_g H) dg$

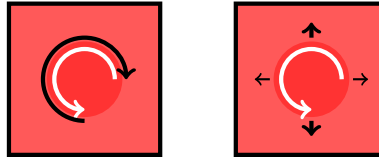


## 1.6 Odd (non-dissipative) response

- White arrow: Driving.
- Black arrow response.



Charge transport



Viscosity



## 1.7 Odd and Isotropic $R$ in 2D

- Odd+isotropic 2-nd rank:  $\propto \underbrace{\begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}}_{\text{Levi-Civita} = i\sigma_y}$
- Viscosity: symmetric tensors  $\mapsto$  symmetric tensors
- Basis for symmetric tensors in 2-D

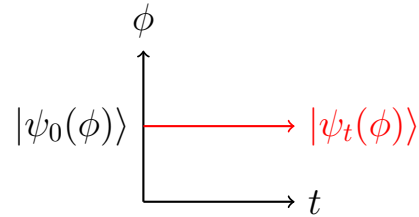
$$\boldsymbol{\sigma} = \underbrace{(\mathbf{1}, \sigma_x, \sigma_z)}_{\text{symmetric tensors}}$$

- Odd isotropic 4-th rank :  $\propto (\boldsymbol{\sigma}_1 \otimes \boldsymbol{\sigma}_3 - \boldsymbol{\sigma}_3 \otimes \boldsymbol{\sigma}_1)$

**Exercise 1.** *Show that in 3 D: Isotropy + no dissipation implies vanishing transport.*

## 1.8 Time dependent Feynman Hellman

- $i\partial_t |\psi\rangle = H |\psi\rangle$
- $\underbrace{\langle \psi_t | \partial_\phi H | \psi_t \rangle}_{\text{virtual work}} = i d_t \left( \underbrace{\langle \psi_t | \partial_\phi \psi_t \rangle}_{\text{geometric}} \right)$



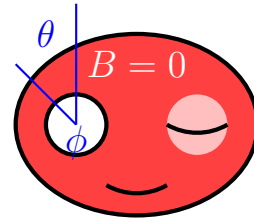
- Proof:

$$\begin{aligned}
 \langle \psi | \partial_\phi H | \psi \rangle &= \partial_\phi \langle \psi | H | \psi \rangle - \langle \partial_\phi \psi | H | \psi \rangle - \langle \psi | H | \partial_\phi \psi \rangle \\
 &= \partial_\phi \langle \psi | \underbrace{H}_{id_t} | \psi \rangle - \langle \partial_\phi \psi | \underbrace{H}_{id_t} | \psi \rangle - \underbrace{\langle \psi | H | \partial_\phi \psi \rangle}_{\langle H \psi |} \\
 &= i\partial_\phi \langle \psi | d_t \psi \rangle - i\langle \partial_\phi \psi | d_t \psi \rangle + i\langle d_t \psi | \partial_\phi \psi \rangle \\
 &= d_t (i\langle \psi | \partial_\phi \psi \rangle)
 \end{aligned}$$

## 2 Topology of control space

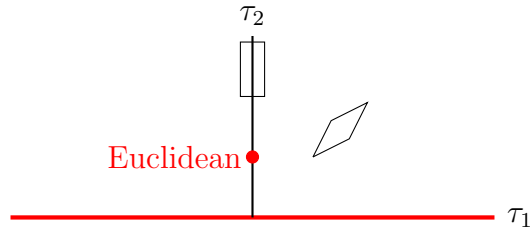
### 2.1 Flux torus

- Flat:  $B = 0$  on red region
- Flux  $\phi$  through hole:  $\oint A = \phi$
- Aharonov-Bohm  $H(\phi) = UH(\phi + 2\pi)U^*$
- Proof:
  - Take  $U = e^{i\theta}$
  - $-iU^*dU = -d\log U$
  - $\oint A \mapsto \phi + 2\pi$
- $\phi$  is an angle: Flux space is a torus



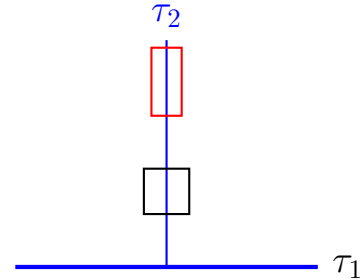
## 2.2 Flat metrics with $\det g = 1$

- $g(\tau) = \frac{1}{\tau_2} \begin{pmatrix} 1 & \tau_1 \\ \tau_1 & |\tau|^2 \end{pmatrix}$ ,  $\tau = \tau_1 + i\tau_2$   $\tau_2 > 0$
- $\det g = \frac{|\tau|^2 - \tau_1^2}{\tau_2^2} = 1$
- $g(\tau)$  flat
  - curvature=0
  - No holonomy of parallel transport.



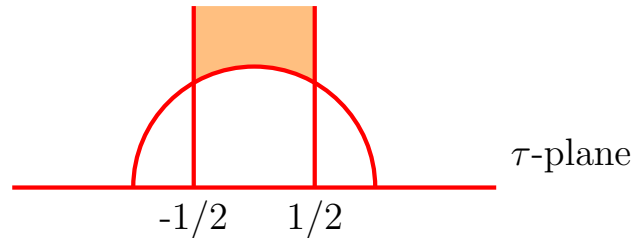
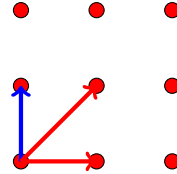
### 2.3 The metric on the space of metrics (Moduli)

- $SL(2, \mathbb{R}) : L = \begin{pmatrix} a & b \\ c & d \end{pmatrix}, \quad \det L = 1, \quad a, b, c, d \in \mathbb{R}$
- $g = (SL)^t \mathbf{1} (SL) = (SL)^t \cancel{SO(2)^t} \mathbf{1} \cancel{SO(2)} (SL)$
- $\tau \leftrightarrow SL(2, \mathbb{R})/SO(2)$
- $g'(\tau') = g(\tau), \quad g' = L^t g L, \quad \tau' = \frac{a\tau+b}{c\tau+d}$
- Haar measure on upper half-plane  $\frac{d\tau_1 d\tau_2}{\tau_2^2}$



## 2.4 The fundamental domain of $SL(2, \mathbb{Z})$

- Torus:  $\mathbb{R}^2/\mathbb{Z}^2$
- $L \in SL(2, \mathbb{Z}) : \mathbb{Z}^2 \mapsto \mathbb{Z}^2$
- $SL(2, \mathbb{R})/SL(2, \mathbb{Z})$ :
- Fundamental domain
  - 2D
  - Finite area
  - Two conic points
  - 1 cusp

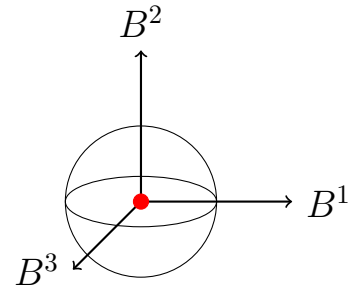
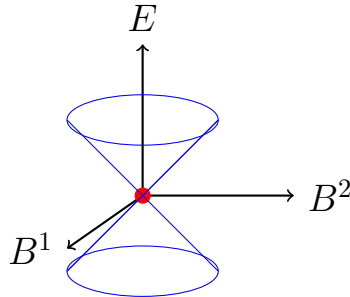


## 2.5 Gapped Hamiltonians

- Gap condition: Endows (Euclidean) control space with interesting topology

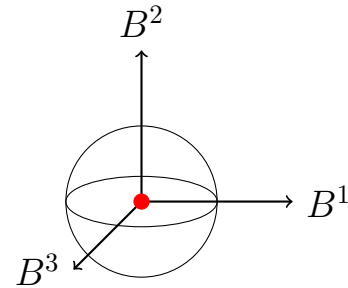
**Example 2.1.**

$$H(\mathbf{B}) = \mathbf{B} \cdot \boldsymbol{\sigma}, \quad \mathbf{B} \in \underbrace{\mathbb{R}^3/0}_{S^2}$$



## 2.6 Wigner von Neuman

- Simple=Non-degenerate
- $\dim(\text{Hermitian } n \times n) = n^2$
- $\dim(\text{Hermitian, 2-fold degenerate}) = n^2 - 3$





## 2.7 Proof

- Simple:  $H = U(n) \cancel{U(1)^{\otimes n}} \underbrace{D}_{\text{diagonal}} \cancel{U^*(1)^{\otimes n}} U^*(n)$

- $H \sim \underbrace{\text{cone}(\lambda_1 < \lambda_2 \cdots < \lambda_n)}_{\text{eigenvalues}} \otimes \underbrace{U(n)/U(1)^n}_{\text{frame}}$

- 1-eigenvalue crossing

$$H = U(n) \cancel{U(1)^{\otimes n-2}} \oplus \cancel{U(2)} \underbrace{(D_{n-2} \oplus \mathbf{1}_2)}_{\text{diagonal } U^*(1)^{\otimes n-2}} \oplus \cancel{U(2)} U^*(n)$$

$$H \sim \text{cone}(\lambda_1 = \lambda_2 \cdots < \lambda_n) \otimes U(n) / \underbrace{U(2)}_{\mathbf{1}_2 V = V \mathbf{1}_2} \otimes U(1)^{n-2}$$

- Co-dimension:  $\underbrace{1}_{\text{cone}} + \underbrace{4}_{U(2)} - \underbrace{2}_{U(1)} = 3$

**Exercise 2.** *co-dim (real symmetric) = 2*

## 2.8 Homotopy of simple matrices

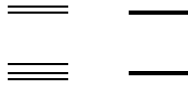
- ( simple hermitian  $n \times n$  )  $\underbrace{\sim}_{\text{homotopic}} U(n)/U(1)^n$
- Deform  $(\lambda_1 < \lambda_2 \dots \lambda_n)$  to  $(1, 2, \dots, n)$ .

$$\begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array} \quad \begin{array}{c} \text{---} \\ \text{---} \\ \text{---} \end{array}$$

- (simple  $2 \times 2$ )  $\sim S^2$

$$\underbrace{\pi_{0,1}(S^2) = 0, \quad \pi_{2,3}(S^2) = \mathbb{Z}, \quad \pi_{5,6}(S^2) = \mathbb{Z}_2}_{\text{rich topology}}$$

## 2.9 Homotopy of gapped matrices

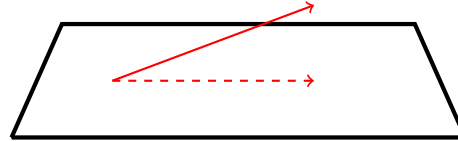


- Gapped matrices homotopic to:

$$U(n + m)/U(n) \otimes U(m)$$

### 3 Geometry of controlled projections

3.1 Projections:  $P^2 = P$

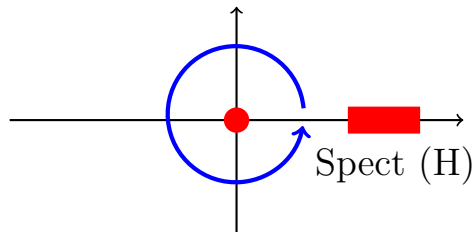


- Pure quantum state

$$P = |\psi\rangle\langle\psi|, \quad \text{Rank}P = 1$$

- Spectral (aka Riesz) projection

$$P = -\frac{1}{2\pi i} \oint R(z)dz, \quad R(z) = \frac{1}{H - z}$$

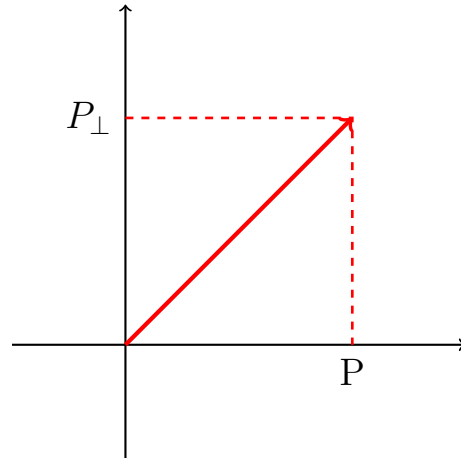


**Exercise 3.** Use

$$(z - z')R(z)R(z') = R(z') - R(z)$$

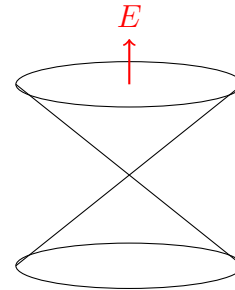
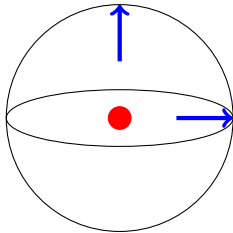
to show that the Riesz integral is a projection.

- Complementary projection:
- $P_{\perp} = \mathbb{1} - P$
- $P_{\perp}^2 = P_{\perp}, \quad PP_{\perp} = P_{\perp}P = 0$



### 3.2 Smooth Families of Projections

- $H(\mathbf{B}) = \mathbf{B} \cdot \boldsymbol{\sigma}$
- Projection on ground state  $P(\mathbf{B}) = \frac{1 - \hat{\mathbf{B}} \cdot \boldsymbol{\sigma}}{2}$
- $(\hat{\mathbf{B}} \cdot \boldsymbol{\sigma})^2 = \mathbb{1} \implies P^2 = P$
- Smooth away from  $\mathbf{B} = 0$



### 3.3 $dP$

- Truly useful:  $P(dP)P = 0$
- Proof:

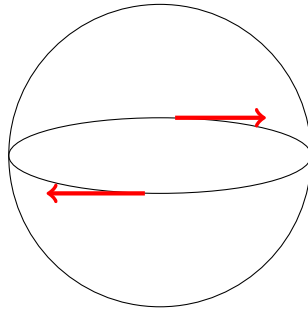
$$P^2 = P \implies PdP + dPP = dP \implies \cancel{P^2dP} + PdPP = \cancel{PdP}$$

- Corollary:

$$PdP = P(dP)P_{\perp} = (dP)P_{\perp}$$

### 3.4 Parallel transport

- Parallel transport = No motion in Range  $P$
- Parallel transport geometry surface
  - Preserves length
  - Preserves angle with geodesic





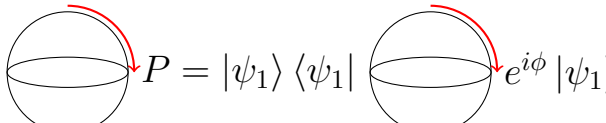
### 3.5 Motion of vectors in Range $P$

- $|\psi\rangle = P |\psi\rangle$
- $d|\psi\rangle = (dP) |\psi\rangle + Pd|\psi\rangle$
- $P_{\perp}d|\psi\rangle = (dP) |\psi\rangle$
- Need an equation for  $Pd|\psi\rangle$
- Parallel transport is the choice:  $Pd|\psi\rangle = 0$

### 3.6 The “geodesic” equation

- The constraint to be in  $P$ :  $P_{\perp} d|\psi\rangle = (dP)|\psi\rangle$
- Together with Parallel transport :  $Pd|\psi\rangle = 0$
- Give the “Geodesic” equation:

$$\begin{aligned} d|\psi\rangle &= (dP)|\psi\rangle \\ &= (dP)P|\psi\rangle \\ &= [dP, P]|\psi\rangle \end{aligned}$$

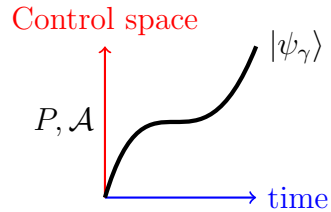
$$P_0 = |\psi_0\rangle\langle\psi_0| \qquad |\psi_0\rangle$$


$$P = |\psi_1\rangle\langle\psi_1| \qquad e^{i\phi} |\psi_1\rangle$$

### 3.7 Connections for projections

$$\underbrace{\mathcal{A} = -i[dP, P]}_{\text{Connection}}, \quad \underbrace{D = d - i\mathcal{A}}_{\text{Covariant derivative}}, \quad \underbrace{D|\psi\rangle = 0}_{\text{Parallel transport}}$$

- $P, \mathcal{A}$ : functions on control space.
- $|\psi_\gamma\rangle$ , solution of ODE: function on space of paths  $\gamma$ .



### 3.8 Unitary intertwining evolutions

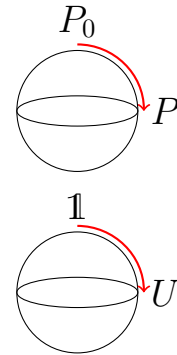
- $\underbrace{PU = UP_0}_{\text{intertwining}}$
- Generator  $\mathcal{A} = i(dU)U^*$
- The equation for the generator

$$\underbrace{dP = i[\mathcal{A}, P]}_{\text{commutator equation}}$$

- Proof

$$\begin{aligned} P_0 = U^*PU \implies 0 &= (dU^*)PU + U^*(dP)U + U^*PdU \\ &= U((dU^*)PU + U^*(dP)U + U^*PdU)U^* \end{aligned}$$

$$0 = U(dU^*)P + (dP) + P(dU)U^* = i\mathcal{A}P + (dP) - iP\mathcal{A}$$



### 3.9 Solutions of the commutator equation

- If  $\mathcal{A}$  solves  $dP = i[\mathcal{A}, P]$
- Then  $\mathcal{A} + \underbrace{\text{Commutant of } P}_{\text{ambiguity}}$  also a solution
- A special solution is Parallel transport

$$\mathcal{A} = -i[dP, P]$$

*Proof.*

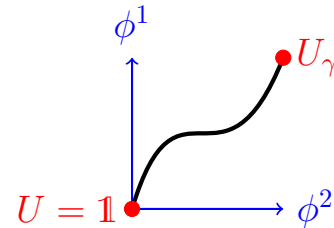
$$i[\mathcal{A}, P] = [[dP, P], P] = (dP)P - 2P(dP)P + PdP = dP$$

□

### 3.10 Evolution in spectral subspaces

- $P$  a spectral projection for  $H$ .
- $\underbrace{P = P^*}_{\text{orthogonal}} \implies \underbrace{\mathcal{A} = \mathcal{A}^*}_{\text{hermitian}} \implies DU = 0, \quad U(0) = \mathbb{1}$
- Generates unitary intertwining evolution  $U_\gamma$ ,
- Parallel transport is re-parametrization invariant

$$dU = -i [dP, P] U + \cancel{(\text{commutant } P)U} dt$$

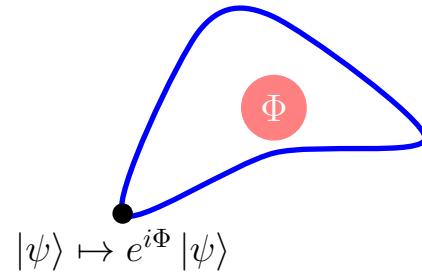


### 3.11 Holonomy

- Closed, anchored path

$$U_\gamma P(\phi_0) = P(\phi_0)U_\gamma$$

- AKA Berry's phase.
- Anchoring point matters (in non-abelian case).



### 3.12 $2\pi$ rotation of a 2-level system

- 2-level system = Qubit
- Projections on equator:  $P(\phi) = \frac{\mathbb{1} + \sigma_x \cos \phi + \sigma_y \sin \phi}{2}$
- Parallel transport = z-rotations  $\underbrace{-i[dP, P]}_{\text{generator of rotations}} = \frac{\sigma_z}{2}$
- Intertwining:  $U(\phi) = e^{i\sigma_z \phi/2} = \mathbb{1} \cos \phi/2 + i\sigma_z \sin \phi/2$
- Holonomy:  $U(2\pi) = -\mathbb{1}$
- $\pi_1(SO(3)) = \mathbb{Z}_2$

